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Stability and Hopf bifurcation in a mathematical model of pluripotent stem cell dynamics

Mostafa Adimy^a, Fabien Crauste^a, Shigui Ruan^{b,*,1}

^aLaboratoire de Mathématiques Appliquées, FRE 2570, Université de Pau et des Pays de l'Adour, Avenue de l'université, 64000 Pau, France

^bDepartment of Mathematics, University of Miami, P.O. Box 249085, Coral Gables, FL 33124-4250, USA

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Abstract

We study a mathematical model describing the dynamics of a pluripotent stem cell population involved in the blood production process in the bone marrow. This model is a differential equation with a time delay. The delay describes the cell cycle duration and is uniformly distributed on an interval. We obtain stability conditions independent of the delay and show that the distributed delay can destabilize the entire system. In particular, it is shown that a Hopf bifurcation can occur. © 2005 Elsevier Ltd. All rights reserved.

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1. Introduction

The blood production process, called hematopoiesis, is one of the major biological phenomena occurring in human body. It takes place in the bone marrow where pluripotent stem cells give birth to mature cells. After ejecting their nuclei, these cells enter the bloodstream and become blood cells.

According to the study of Burns and Tannock [4], the population of pluripotent stem cells can be divided into two distinct groups: quiescent cells and proliferating cells. Mathematical

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^{*} Corresponding author. Tel.: +1 305 284 2312; fax: +1 305 284 2848.

E-mail addresses: mostafa.adimy@univ-pau.fr (M. Adimy), fabien.crauste@univ-pau.fr (F. Crauste), ruan@math.miami.edu (S. Ruan).

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models describing the dynamics of this cell population have been proposed since the end of the 1970s, in particularly by Mackey [9,10]. We refer to the review articles by Haurie et al. [8] and Mackey et al. [12] for further study and more references on this topic. More recently, Pujo-Menjouet et al. [14] and Pujo-Menjouet and Mackey [15] proved the existence of a Hopf bifurcation for the hematopoiesis model proposed in [9]. In all these works, the authors assumed that the proliferating phase duration is constant. Mathematically, this means that the delay in these models is a discrete delay. However, experimental data (see [3]) indicate that cells do not spend the same time in the proliferating phase.

In this paper, taking into account this assumption, we assume that the delay (or proliferating phase duration) is uniformly distributed on an interval. The main objective is to investigate the effect of time delay on the dynamical solutions. It is shown that there exist some critical values of the time delay such that a local Hopf bifurcation occurs at the nontrivial equilibrium.

The paper is organized as follows. In Section 2, we present our model, which is given in Eq. (1). In Section 3, we derive stability conditions for the two equilibria of Eq. (1) which do not depend on the delay. We show the existence of a Hopf bifurcation at the nontrivial equilibrium in Section 4. A brief discussion is given in Section 5.

2. The model

Pluripotent stem cells can be either in a resting phase, also known as G_0 -phase, or in a proliferating phase. In the resting phase, they can die at a constant rate $\delta \ge 0$, which also includes the cellular differentiation, or be introduced in the proliferating phase at a rate β . According to the work of Sachs [16], β is assumed to depend on the resting phase population.

In the proliferating phase, which is in fact the so-called cell cycle, pluripotent stem cells are committed to divide and give birth to two daughter cells at the end of this phase. The two daughter cells enter directly the resting phase and complete the cycle. We assume that proliferating cells divide according to a uniform law *f* on an interval $[\tau_{\min}, \tau]$ with $0 \le \tau_{\min} < \tau < +\infty$. This assumption comes from the fact that, even if only a little is known about phenomena involved in hematopoiesis, there are strong evidences (see [3]) indicating that cells do not divide at the same age. The function *f* is then defined by

$$f(r) = \begin{cases} \frac{1}{\tau - \tau_{\min}} & \text{if } r \in [\tau_{\min}, \tau], \\ 0 & \text{otherwise.} \end{cases}$$

Let x(t) denote the pluripotent stem cell population density (cells/kg) at time $t \ge 0$. It satisfies the nonlinear delay differential equation

$$x'(t) = -(\delta + \beta(x(t)))x(t) + \frac{2}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} \beta(x(t-r))x(t-r) \,\mathrm{d}r.$$
(1)

The first term in the right-hand side of Eq. (1) accounts for the cellular loss due to mortality and cellular differentiation, $\delta x(t)$, and introduction in the cell cycle, $\beta(x(t))x(t)$. The second term is for the division of proliferating cells into two daughter cells during mitosis. Proliferating cells are in fact resting cells introduced in the proliferating phase one generation earlier, so that the quantity $\beta(x(t-r))x(t-r)$ appears with a time delay. The factor 2 is, of course, for the division of each proliferating cell into two daughter cells.

In the following, the rate of reintroduction in the proliferating compartment β is taken to be a monotone and decreasing Hill function, given by

$$\beta(x) = \beta_0 \frac{\theta^n}{\theta^n + x^n} \quad \text{for } x \ge 0.$$

The coefficient $\beta_0 > 0$ is the maximum rate of reintroduction, $\theta \ge 0$ is the G_0 -phase population density for which the rate of re-entry β attains its maximum rate of change with respect to the resting phase population, and $n \ge 0$ describes the sensitivity of β with changes in the population. This function was first used in hematopoiesis models by Mackey [9] in 1978.

In [9,11], Mackey gave values of the above parameters for a normal human body production. These values are

$$\delta = 0.05 \mathrm{d}^{-1}, \quad \beta_0 = 1.77 \mathrm{d}^{-1} \quad \text{and} \quad n = 3.$$
 (2)

The value of θ is usually $\theta = 1.62 \times 10^8$ cells/kg. However, since we shall study the qualitative behavior of the pluripotent stem cells population, the value of θ is not really important and could be normalized without loss of generality.

Now if we consider an initial continuous nonnegative function φ defined on $[-\tau, 0]$, then Eq. (1) has a unique continuous and nonnegative solution $x^{\varphi}(t)$, defined for $t \ge -\tau$, such that

$$x^{\varphi}(s) = \varphi(s) \quad \text{for } s \in [-\tau, 0].$$

This can be obtained by using the results in [7].

Notice that Eq. (1) has at most two equilibria, the trivial equilibrium $x \equiv 0$ and a nontrivial positive equilibrium $x \equiv x^*$. The trivial equilibrium always exists and corresponds to the extinction of the population.

Proposition 2.1. *Eq.* (1) *has a nontrivial positive equilibrium* $x \equiv x^*$ *if and only if*

$$\beta_0 > \delta > 0. \tag{3}$$

In this case, x^* is explicitly given by

$$x^* = \theta \left(\frac{\beta_0}{\delta} - 1\right)^{1/n}.$$

Proof. Let x^* be an equilibrium of Eq. (1). Then x^* satisfies $x^* = 0$ or $\beta(x^*) - \delta = 0$. Consequently, Eq. (1) has a nontrivial equilibrium if and only if the equation $\beta(x^*) = \delta$ has a nontrivial solution. Since the function β is decreasing and positive with $\beta(0) = \beta_0$, Eq. (1) has a nontrivial equilibrium if and only if condition (3) holds.

In the next section, we shall study the stability of the two equilibria of Eq. (1).

3. Stability

Throughout this section, we are interested in the stability of the equilibria of Eq. (1), in particularly the stability of the nontrivial equilibrium $x \equiv x^*$. We start by giving a result on the global stability of the trivial equilibrium of (1).

Theorem 3.1. The trivial equilibrium $x \equiv 0$ of Eq. (1) is globally stable if

 $\beta_0 < \delta$.

Proof. The proof uses a similar technique employed by Adimy and Crauste [1]. It is based on the construction of a Lyapunov functional.

Denote by C^+ the space of all continuous nonnegative functions on $[-\tau, 0]$. Consider the mapping $J : C^+ \to [0, +\infty)$ defined, for $\varphi \in C^+$, by

$$J(\varphi) = \int_0^{\varphi(0)} \beta(r) r \, \mathrm{d}r + \frac{1}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} \int_{-r}^0 (\beta(\varphi(a))\varphi(a))^2 \, \mathrm{d}a \, \mathrm{d}r.$$

Then,

$$\begin{split} \hat{J}(\varphi) &= \dot{\varphi}(0)\beta(\varphi(0))\varphi(0) \\ &+ \frac{1}{\tau - \tau_{\min}} \left[\int_{\tau_{\min}}^{\tau} (\beta(\varphi(0))\varphi(0))^2 - (\beta(\varphi(-r))\varphi(-r))^2 \right] \mathrm{d}r. \end{split}$$

Since

$$\dot{\varphi}(0) = -(\delta + \beta(\varphi(0)))\varphi(0) + \frac{2}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} \beta(\varphi(-r))\varphi(-r) \,\mathrm{d}r,$$

we obtain that

$$\dot{J}(\varphi) = -(\delta + \beta(\varphi(0)))\beta(\varphi(0))\varphi^{2}(0) + \frac{2}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} (\beta(\varphi(0))\varphi(0))^{2} dr - \frac{1}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} (\beta(\varphi(0))\varphi(0) - \beta(\varphi(-r))\varphi(-r))^{2} dr.$$

Hence,

$$\dot{J}(\varphi) \leqslant -(\delta - \beta(\varphi(0)))\beta(\varphi(0))\varphi(0)^2$$

Let α be the function defined, for $x \ge 0$, by

$$\alpha(x) = (\delta - \beta(x))\beta(x)x^2.$$

Assume that $\beta_0 < \delta$. Since β is a decreasing function, it follows that the function $x \mapsto \delta - \beta(x)$ is positive for $x \ge 0$. Hence, α is nonnegative on $[0, +\infty)$ and $\alpha(x) = 0$ if and only if x = 0. Consequently, the mapping *J* is a Lyapunov functional when $\beta_0 < \delta$. We then deduce that the trivial equilibrium of (1) is globally stable. \Box

The result in Theorem 3.1 describes the fact that when $x \equiv 0$ is the only equilibrium of (1), the population is doomed to extinction except when $\beta_0 = \delta$.

Now we focus on the stability of the positive equilibrium $x \equiv x^*$ of Eq. (1). To ensure the existence of the equilibrium $x \equiv x^*$, we assume that condition (3) holds; that is, $\beta_0 > \delta > 0$. We do not expect to obtain conditions for the global stability of $x \equiv x^*$. However, local stability results can be obtained by linearizing equation (1) about x^* . Set

$$\beta^* := \left. \frac{\mathrm{d}}{\mathrm{d}x} (\beta(x)x) \right|_{x=x^*} = \delta\left(1 - n \, \frac{\beta_0 - \delta}{\beta_0} \right). \tag{4}$$

The linearization of Eq. (1) at x^* is

$$x'(t) = -(\delta + \beta^*)x(t) + \frac{2\beta^*}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} x(t-r) \,\mathrm{d}r.$$

The characteristic equation of (1) is given by

$$\Delta(\lambda) := \lambda + \delta + \beta^* - \frac{2\beta^*}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} e^{-\lambda r} dr = 0.$$
⁽⁵⁾

We now state and prove our first result on the stability of $x \equiv x^*$.

Theorem 3.2. Assume that

$$n \, \frac{\beta_0 - \delta}{\beta_0} \leqslant 1. \tag{6}$$

Then the nontrivial equilibrium $x \equiv x^*$ of Eq. (1) is locally asymptotically stable.

Proof. Assume that (6) holds, which is equivalent, by using (4), to $\beta^* \ge 0$.

We first assume that $\Delta(\lambda)$, given by (5), is a real function. Then, $\Delta(\lambda)$ is continuously differentiable and its first derivative is given by

$$\frac{d\Delta}{d\lambda}(\lambda) = 1 + \frac{2\beta^*}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} r e^{-\lambda r} dr.$$
(7)

One can see that $d\Delta/d\lambda$ is positive for $\lambda \in \mathbb{R}$ as soon as $\beta^* \ge 0$. Moreover,

$$\lim_{\lambda \to -\infty} \Delta(\lambda) = -\infty \quad \text{and} \quad \lim_{\lambda \to +\infty} \Delta(\lambda) = +\infty.$$

Consequently, $\Delta(\lambda)$ has a unique real root λ_0 . Since

$$\Delta(0) = \delta - \beta^* = n\delta\left(1 - \frac{\delta}{\beta_0}\right) > 0,$$

we deduce that $\lambda_0 < 0$.

Now, we show that if λ is a characteristic root of Eq. (5), then Re(λ) $\leq \lambda_0$. By contradiction, we assume that there exists a characteristic root $\lambda = \mu + i\omega$ of Eq. (5) such that $\mu > \lambda_0$. By considering the real part of $\Delta(\lambda)$, we obtain that

$$\mu + \delta + \beta^* - \frac{2\beta^*}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} e^{-\mu r} \cos(\omega r) dr = 0.$$

Consequently,

$$\mu - \lambda_0 = \frac{2\beta^*}{\tau - \tau_{\min}} \int_{\tau_{\min}}^{\tau} (e^{-\mu r} \cos(\omega r) - e^{-\lambda_0 r}) dr < 0.$$

This yields a contradiction. We conclude that every characteristic root λ of (5) is such that $\operatorname{Re}(\lambda) \leq \lambda_0$. Hence, all characteristic roots of (5) have negative real parts and the equilibrium $x \equiv x^*$ is locally asymptotically stable. \Box

When $\beta^* < 0$, that is, when

$$1 < n \, \frac{\beta_0 - \delta}{\beta_0},$$

the stability cannot occur for all values of τ_{\min} and τ . In particularly, we shall show that a Hopf bifurcation can occur (see Theorem 4.5). However, we can still have the stability of the nontrivial equilibrium $x \equiv x^*$ for values of n, β_0 and δ if $n(\beta_0 - \delta)/\beta_0$ is not too large. This will be considered in the next theorem.

To present the results, without loss of generality, we assume that

$$\tau_{\min} = 0.$$

We want to point out that the results we are going to show remain true when $\tau_{\min} > 0$, but the proof is more complicated.

Define a function *K*, for $x \ge 0$, by

$$K(x) = \frac{\sin(x)}{x} \tag{8}$$

and let x_1 be the unique solution of the equation

$$x_1 = \tan(x_1), \quad x_1 \in \left(\pi, \frac{3\pi}{2}\right).$$

Set

$$u_0 := \cos(x_1) \in (-1, 0).$$

Then $K'(x_1) = 0$ and

$$u_0 = K(x_1) = \min_{x \ge 0} K(x).$$
(9)

We have the following local stability theorem.

Theorem 3.3. Assume that

$$1 < n \,\frac{\beta_0 - \delta}{\beta_0} < \frac{2(1 - u_0)}{1 - 2u_0}.\tag{10}$$

Then the nontrivial equilibrium $x \equiv x^*$ of Eq. (1) is locally asymptotically stable.

Proof. Let us assume that (10) holds. Then $\beta^* < 0$, $\delta + \beta^* > 0$ and

$$\frac{\delta + \beta^*}{2\beta^*} < u_0. \tag{11}$$

By contradiction, assume that there exists a characteristic root $\lambda = \mu + i\omega$ of (5) with $\mu > 0$. Then,

$$\mu = -(\delta + \beta^*) + \frac{2\beta^*}{\omega\tau} \int_0^{\omega\tau} e^{-(\mu/\omega)r} \cos(r) dr$$

and

$$\omega = -\frac{2\beta^*}{\omega\tau} \int_0^{\omega\tau} e^{-(\mu/\omega)r} \sin(r) dr.$$

Integrating by parts, we obtain that

$$2\mu = -(\delta + \beta^*) + 2\beta^* \mathrm{e}^{-\mu\tau} K(\omega\tau).$$

Consequently,

$$\mu < -(\delta + \beta^*) + 2\beta^* \mathrm{e}^{-\mu\tau} K(\omega\tau).$$

If $\omega \tau$ is such that $\sin(\omega \tau) \ge 0$, then $2\beta^* K(\omega \tau) < 0$ and

 $\mu < -(\delta + \beta^*) \leqslant 0.$

So we obtain a contradiction.

Similarly, if $\omega \tau$ is such that $\sin(\omega \tau) < 0$, then, from (9) and (11), we deduce that

$$\frac{\delta + \beta^*}{2\beta^*} \leqslant K(\omega\tau).$$

It implies that

$$\delta + \beta^* \ge 2\beta^* K(\omega\tau) > 2\beta^* \mathrm{e}^{-\mu\tau} K(\omega\tau).$$

Therefore,

$$\mu < -(\delta + \beta^*) + 2\beta^* K(\omega \tau) \leq 0.$$

Again we obtain a contradiction. Hence, all characteristic roots λ of (5) are such that $\operatorname{Re}(\lambda) \leq 0$.

Now, we assume that (5) has a purely imaginary characteristic root $\lambda = i\omega$. Then ω and τ satisfy

$$K(\omega\tau) = \frac{\delta + \beta^*}{2\beta^*}.$$
(12)

Using (9) and (11), we obtain a contradiction. Consequently, (12) has no solution and Eq. (5) does not have purely imaginary roots. We conclude that all characteristic roots of (5) have negative real parts and $x \equiv x^*$ is locally asymptotically stable. \Box

From Theorems 3.2 and 3.3, it follows that the nontrivial equilibrium $x \equiv x^*$ of Eq. (1) is locally asymptotically stable when

$$0 \leq n \, \frac{\beta_0 - \delta}{\beta_0} < \frac{2(1 - u_0)}{1 - 2u_0}.\tag{13}$$

We are going to show that as soon as condition (13) does not hold, then the equilibrium can be destabilized. In the next section, we shall show that if condition (11) does not hold, then a Hopf bifurcation indeed occurs at $x \equiv x^*$.

4. Hopf bifurcation

In this section, we are going to show that the nontrivial equilibrium $x \equiv x^*$ of Eq. (1) can be destabilized *via* a Hopf bifurcation. The time delay τ will be used as a bifurcation parameter. This result is obtained in Theorem 4.5.

Recall that the nontrivial equilibrium $x \equiv x^*$ of Eq. (1) exists if and only if $\beta_0 > \delta > 0$. In the following, without loss of generality, we assume that

$$\tau_{\min}=0.$$

Again the results still hold when $\tau_{\min} > 0$, but the proof is easier to understand when $\tau_{\min} = 0$.

We look for purely imaginary roots of $\Delta(\lambda)$. Of course, we assume that $\beta^* < 0$, otherwise $x \equiv x^*$ is locally asymptotically stable. Let $\lambda = i\omega$, with $\omega \in \mathbb{R}$, be a purely imaginary characteristic root of Eq. (5). Then, τ and ω satisfy the following system:

$$\begin{cases} \delta + \beta^* (1 - 2C(\tau, \omega)) = 0, \\ \omega + 2\beta^* S(\tau, \omega) = 0, \end{cases}$$
(14)

where

$$C(\tau, \omega) = \frac{1}{\tau} \int_0^\tau \cos(\omega r) \, \mathrm{d}r, \quad S(\tau, \omega) = \frac{1}{\tau} \int_0^\tau \sin(\omega r) \, \mathrm{d}r.$$

First, one can see that $\omega = 0$ cannot be a solution of (14). Otherwise $\delta = \beta^* < 0$. Moreover, if ω is a solution of system (14), then $-\omega$ is also a solution of (14). Hence, we only look for positive solutions ω .

One can check that $C(\tau, \omega)$ and $S(\tau, \omega)$ are given, for $\tau > 0$ and $\omega > 0$, by

$$C(\tau, \omega) = \frac{\sin(\omega\tau)}{\omega\tau} = K(\omega\tau), \quad S(\tau, \omega) = \frac{1 - \cos(\omega\tau)}{\omega\tau},$$

where the function K is defined by (8). Consequently, system (14) can be rewritten as

$$K(\omega\tau) = \frac{\delta + \beta^*}{2\beta^*},\tag{15}$$

$$\frac{\cos(\omega\tau) - 1}{(\omega\tau)^2} = \frac{1}{2\beta^*\tau}.$$
(16)

Consider the sequence

$$\{x_k\}_{k \in \mathbb{N}} := \{x \ge 0; \ x = \tan(x)\}$$
(17)

with

 $0 = x_0 < x_1 < \cdots < x_k < \cdots$

In fact, one can check that

$$\{x_k\}_{k \in \mathbb{N}} = \{x \ge 0; K'(x) = 0\}.$$

Moreover, for all $k \in \mathbb{N}^*$,

$$x_k \in \left(k\pi, k\pi + \frac{\pi}{2}\right).$$

Define two sequences $\{u_k\}$ and $\{v_k\}$, for $k \in \mathbb{N}$, by

 $u_k := \cos(x_{2k+1}) < 0, \quad v_k := \cos(x_{2k}) > 0.$

Using the definition of x_k , one can see that

 $u_k = K(x_{2k+1})$ and $v_k = K(x_{2k})$.

Thus, the sequence $\{u_k\}_{k\in\mathbb{N}}$ is increasing with $-1 < u_k < 0$ and the sequence $\{v_k\}_{k\in\mathbb{N}}$ is decreasing with $v_0 = 1$ and $0 < v_k < 1/2$ for $k \ge 1$ (see Fig. 1).

Moreover,

$$\lim_{k \to +\infty} u_k = \lim_{k \to +\infty} v_k = 0.$$

Furthermore, one can check that, as soon as $\beta^* < 0$,

$$\frac{\delta + \beta^*}{2\beta^*} < 1 = v_0.$$

Finally, define a function *h*, for $x \in [-1, 1/2)$, by

$$h(x) = \frac{2(1-x)}{1-2x}$$

and set

 $h(v_0) = +\infty.$

We have the following results about the properties of the function h.

Lemma 4.1. Suppose that

$$h(u_0) \leq n \frac{\beta_0 - \delta}{\beta_0} \quad and \quad \delta + \beta^* \neq 0.$$

(i) If $\delta + \beta^* > 0$, then there exists $k \in \mathbb{N}$ such that

$$h(u_k) \leqslant n \, \frac{\beta_0 - \delta}{\beta_0} < h(u_{k+1}).$$



Fig. 1. The graph of K(x).

(ii) If
$$\delta + \beta^* < 0$$
, then there exists $k \in \mathbb{N}$ such that

$$h(v_{k+1}) \leq n \frac{\beta_0 - \delta}{\beta_0} < h(v_k)$$

Proof. Since the function *h* is increasing on the interval [-1, 1/2), we can see that

$$h(u_k) \leqslant n \, \frac{\beta_0 - \delta}{\beta_0} < h(u_{k+1})$$

is equivalent to

$$u_k \leqslant \frac{\delta + \beta^*}{2\beta^*} < u_{k+1}.$$

Also,

$$h(v_{k+1}) \leqslant n \, \frac{\beta_0 - \delta}{\beta_0} < h(v_k)$$

is equivalent to

$$v_{k+1} \leqslant \frac{\delta + \beta^*}{2\beta^*} < v_k.$$

The lemma now follows. \Box

Proposition 4.2. If

$$h(u_k) < n \, \frac{\beta_0 - \delta}{\beta_0} < h(u_{k+1}), \quad k \in \mathbb{N},$$

then system (15)–(16) *has exactly* 2(k + 1) *solutions* $(\tau_{1,1}, \omega_{1,1}), \ldots, (\tau_{k+1,1}, \omega_{k+1,1})$ *and* $(\tau_{1,2}, \omega_{1,2}), \ldots, (\tau_{k+1,2}, \omega_{k+1,2})$ *with*

$$\begin{cases} \omega_{l,1}\tau_{l,1} \in ((2l-1)\pi, x_{2l-1}) & \text{for } l = 1, \dots, k+1, \\ \omega_{l,2}\tau_{l,2} \in (x_{2l-1}2l\pi) & \text{for } l = 1, \dots, k+1 \end{cases}$$

and

$$0 < \tau_{1,1} < \cdots < \tau_{k+1,1} < \tau_{k+1,2} < \cdots < \tau_{1,2}.$$

(ii) If

$$n \, \frac{\beta_0 - \delta}{\beta_0} = h(u_k), \quad k \in \mathbb{N},$$

then system (15)–(16) has exactly 2k + 1 solutions $(\tau_{1,1}, \omega_{1,1}), \ldots, (\tau_{k+1,1}, \omega_{k+1,1})$ and $(\tau_{1,2}, \omega_{1,2}), \ldots, (\tau_{k,2}, \omega_{k,2})$ with

$$\begin{cases} \omega_{l,1}\tau_{l,1} \in ((2l-1)\pi, x_{2l-1}) & \text{for } l = 1, \dots, k, \\ \omega_{l,2}\tau_{l,2} \in (x_{2l-1}, 2l\pi) & \text{for } l = 1, \dots, k, \\ \omega_{k+1,1}\tau_{k+1,1} = x_{2k+1} \end{cases}$$

and

$$0 < \tau_{1,1} < \cdots < \tau_{k+1,1} < \tau_{k,2} < \cdots < \tau_{1,2}.$$

(iii) If

$$h(v_{k+1}) < n \, \frac{\beta_0 - \delta}{\beta_0} < h(v_k), \quad k \in \mathbb{N}^*,$$

then system (15)–(16) has exactly 2k + 1 solutions $(\tau_{1,1}, \omega_{1,1}), \ldots, (\tau_{k+1,1}, \omega_{k+1,1})$ and $(\tau_{1,2}, \omega_{1,2}), \ldots, (\tau_{k,2}, \omega_{k,2})$ with

$$\begin{cases} \omega_{1,1}\tau_{l,1} \in (\pi/2, \pi), \\ \omega_{l,1}\tau_{l,1} \in (x_{l+1}, (l+2)\pi) & \text{for } l = 2, \dots, k+1, \\ \omega_{l,2}\tau_{l,2} \in ((l+1)\pi, x_{l+1}) & \text{for } l = 1, \dots, k \end{cases}$$

and

$$0 < \tau_{1,1} < \cdots < \tau_{k+1,1} < \tau_{k,2} < \cdots < \tau_{1,2}.$$

(iv) If

$$n \, \frac{\beta_0 - \delta}{\beta_0} = h(v_k), \quad k \in \mathbb{N}^*,$$

then system (15)–(16) has exactly 2k solutions $(\tau_{1,1}, \omega_{1,1}), \ldots, (\tau_{k,1}, \omega_{k,1})$ and $(\tau_{1,2}, \omega_{1,2}), \ldots, (\tau_{k,2}, \omega_{k,2})$, with

$$\begin{array}{ll} \omega_{l,1}\tau_{l,1} \in (\pi/2,\pi), \\ \omega_{l,1}\tau_{l,1} \in (x_l,(l+1)\pi) & for \ l=2,\ldots,k, \\ \omega_{l,2}\tau_{l,2} \in ((l+1)\pi,x_{l+1}) & for \ l=1,\ldots,k-1, \\ \omega_{k,2}\tau_{k,2} = x_{2k} \end{array}$$

and

$$0 < \tau_{1,1} < \cdots < \tau_{k,1} < \tau_{k,2} < \cdots < \tau_{1,2}$$

(v) *If*

$$h(v_1) < n \, \frac{\beta_0 - \delta}{\beta_0} < h(v_0),$$

then system (15)–(16) has a unique solution (τ_1, ω_1) such that $\tau_1 > 0$ and

$$\omega_1 \tau_1 \in (0, \pi).$$

Proof. We only prove (i) when k = 0. The other cases can be deduced similarly. Assume that

$$h(u_0) < n \frac{\beta_0 - \delta}{\beta_0} < h(u_1).$$

This is equivalent to

$$u_0 < \frac{\delta + \beta^*}{2\beta^*} < u_1.$$

The function *K* is strictly negative and decreasing on (π, x_1) with $K(y) \in (u_0, 0)$ (see Fig. 1). So the equation

$$K(y) = \frac{\delta + \beta^*}{2\beta^*}$$

has a unique solution y_1 on the interval (π, x_1) . Set

$$\tau_{1,1} = \frac{(y_1)^2}{2\beta^*(\cos(y_1) - 1)}$$

and $\omega_{1,1} = y_1/\tau_{1,1}$. Then, $(\tau_{1,1}, \omega_{1,1})$ is the unique solution of system (15)–(16) satisfying $\omega_{1,1}\tau_{1,1} \in (\pi, x_1)$.

Moreover, the function *K* is strictly negative and increasing on $(x_1, 2\pi)$ with $K(y) \in (u_0, 0)$, so the equation $K(y) = (\delta + \beta^*)/2\beta^*$ has a unique solution y_2 on the interval $(x_1, 2\pi)$. Set

$$\tau_{1,2} = \frac{(y_2)^2}{2\beta^*(\cos(y_2) - 1)}$$

and $\omega_{1,2} = y_2/\tau_{1,2}$. Then, $(\tau_{1,2}, \omega_{1,2})$ is the unique solution of system (15)–(16) which satisfies $\omega_{1,2}\tau_{1,2} \in (x_1, 2\pi)$.

Furthermore, the function *K* is nonnegative on $[0, \pi]$ and

$$u_1 = K(x_3) = \min_{x \ge 2\pi} K(x).$$

Therefore, system (15)–(16) has two solutions, $(\tau_{1,1}, \omega_{1,1})$ and $(\tau_{1,2}, \omega_{1,2})$.

Finally, using the fact that $\cos(y_1) \leq \cos(y_2)$, we obtain that $\tau_{1,1} < \tau_{1,2}$. This completes the proof. \Box

Lemma 4.1 and Proposition 4.2 give conditions for the existence of pairs of purely imaginary roots of Eq. (5). In the next proposition, we study the properties of the purely imaginary roots of (5).

Proposition 4.3. Assume that there exists a $\tau_c > 0$ such that Eq. (5) has a pair of purely imaginary roots $\pm i\omega_c$ for $\tau = \tau_c$ with $\omega_c > 0$. If

$$\omega_c \tau_c \neq x_k \quad for \ all \ k \in \mathbb{N},$$

where the sequence $\{x_k\}_{k\in\mathbb{N}}$ is defined by (17), then $\pm i\omega_c$ are simple roots such that

$$\begin{cases} \left. \frac{\mathrm{d}\operatorname{Re}(\lambda)}{\mathrm{d}\tau} \right|_{\tau=\tau_c} > 0 \quad if \ \omega_c \tau_c \in (x_{2k}, x_{2k+1}), \\ \left. \frac{\mathrm{d}\operatorname{Re}(\lambda)}{\mathrm{d}\tau} \right|_{\tau=\tau_c} < 0 \quad if \ \omega_c \tau_c \in (x_{2k+1}, x_{2k+2}), \quad k \in \mathbb{N}. \end{cases}$$
(18)

Proof. Assume that there exists a $\tau_c > 0$ such that Eq. (5) has a pair of purely imaginary roots $\pm i\omega_c$ for $\tau = \tau_c$ with $\omega_c > 0$. Then, $\omega_c \tau_c$ satisfies system (15)–(16).

Assume that $\omega_c \tau_c \neq x_k$ for all $k \in \mathbb{N}$. Let us show that $\pm i\omega_c$ are simple characteristic roots of (5). Using (7), one can see that $\pm i\omega_c$ are simple roots of (5) if

$$1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \neq 0 \quad \text{or} \quad \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \neq 0.$$

We will show that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \neq 0.$$

A simple computation shows that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c) = \frac{g(\omega_c \tau_c)}{\omega_c^2 \tau_c},$$

where the function g is defined by

$$g(x) = x \cos(x) - \sin(x)$$
 for $x \ge 0$.

One can check that g(x) = 0 if and only if there exists a $k_0 \in \mathbb{N}$ such that $x = x_{k_0}$. Moreover,

g(x) > 0 if and only if $x \in (x_{2k+1}, x_{2k+2}), k \in \mathbb{N}$.

This yields that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c) < 0 \quad \text{if } \omega_c \tau_c \in (x_{2k}, x_{2k+1})$$
(19)

and

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c) > 0 \quad \text{if } \omega_c \tau_c \in (x_{2k+1}, x_{2k+2}).$$
(20)

Hence, $\pm i\omega_c$ are simple characteristic roots of (5).

Let $\lambda(\tau) = \mu(\tau) + i\omega(\tau)$ be a characteristic root of (5) such that $\lambda(\tau_c) = \pm i\omega_c$. By separating the real and imaginary parts, we obtain that

$$\begin{cases} \mu(\tau) + \delta + \beta^* - \frac{2\beta^*}{\tau} \int_0^\tau e^{-\mu(\tau)r} \cos(\omega(\tau)r) \, dr = 0, \\ \omega(\tau) + \frac{2\beta^*}{\tau} \int_0^\tau e^{-\mu(\tau)r} \sin(\omega(\tau)r) \, dr = 0. \end{cases}$$

We denote by $\mu'(\tau)$ (respectively $\omega'(\tau)$) the first derivative of $\mu(\tau)$ (respectively $\omega(\tau)$) with respect to τ . For $\tau = \tau_c$, we obtain that

$$\mu'(\tau_c) \left[1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \right]$$

= $2\beta^* \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \omega'(\tau_c) + \frac{2\beta^*}{\tau_c}(\cos(\omega_c \tau_c) - C(\tau_c, \omega_c))$ (21)

and

$$\omega'(\tau_c) \left[1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \right]$$

= $-2\beta^* \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \mu'(\tau_c) + \frac{2\beta^*}{\tau_c}(S(\tau_c, \omega_c) - \sin(\omega_c \tau_c)).$ (22)

We consider two cases. First, assume that

$$1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) = 0.$$
⁽²³⁾

One can verify that

$$1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) = 2 + (\delta + \beta^*)\tau_c$$

Consequently, (23) is equivalent to

$$\tau_c = -\frac{2}{\delta + \beta^*}.\tag{24}$$

Then, it follows from Eq. (22) that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \mu'(\tau_c) = \frac{S(\tau_c, \omega_c) - \sin(\omega_c \tau_c)}{\tau_c}.$$

Moreover, by using (15) and (16), we have

$$\frac{S(\tau_c, \omega_c) - \sin(\omega_c \tau_c)}{\tau_c} = \frac{1 - (\cos(\omega_c \tau_c) + \omega_c \tau_c \sin(\omega_c \tau_c))}{\omega_c \tau_c^2}$$
$$= -\frac{\delta + \beta^*}{4\beta^*} \omega_c.$$

Hence, (24) implies that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c)\mu'(\tau_c) = \frac{\omega_c}{2\beta^*\tau_c} < 0.$$

Since $(\partial C/\partial \omega)(\tau_c, \omega_c) \neq 0$, we have

$$\mu'(\tau_c) \neq 0.$$

Furthermore, the sign of $\mu'(\tau_c)$ is the same as the sign of $-(\partial C/\partial \omega)(\tau_c, \omega_c)$.

We now assume that

$$1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \neq 0.$$

Then, by using (21) and (22), we obtain that $\mu'(\tau_c)$ satisfies

$$\begin{split} \mu'(\tau_c) & \left[\left(1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \right)^2 + \left(2\beta^* \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \right)^2 \right] \\ &= \frac{2\beta^*}{\tau_c} \left[2\beta^* \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) (S(\tau_c, \omega_c) - \sin(\omega_c \tau_c)) \right. \\ & \left. + \left(1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \right) (\cos(\omega_c \tau_c) - C(\tau_c, \omega_c)) \right]. \end{split}$$

Using the definitions of C and S, one can check that

$$\frac{\partial C}{\partial \omega}(\tau_c, \omega_c)(S(\tau_c, \omega_c) - \sin(\omega_c \tau_c)) + \frac{\partial S}{\partial \omega}(\tau_c, \omega_c)(\cos(\omega_c \tau_c) - C(\tau_c, \omega_c)) = 0.$$

Hence,

$$\mu'(\tau_c) \left[\left(1 + 2\beta^* \frac{\partial S}{\partial \omega}(\tau_c, \omega_c) \right)^2 + \left(2\beta^* \frac{\partial C}{\partial \omega}(\tau_c, \omega_c) \right)^2 \right]$$
$$= \frac{2\beta^*}{\tau_c} (\cos(\omega_c \tau_c) - C(\tau_c, \omega_c)).$$

Notice that

$$\frac{\cos(\omega_c \tau_c) - C(\tau_c, \omega_c)}{\tau_c} = \frac{g(\omega_c \tau_c)}{\omega_c \tau_c^2} = \frac{\omega_c}{\tau_c} \frac{\partial C}{\partial \omega} (\tau_c, \omega_c) \neq 0.$$

Since $1 + 2\beta^* (\partial S / \partial \omega)(\tau_c, \omega_c) \neq 0$, it follows that

$$\left(1+2\beta^*\frac{\partial S}{\partial\omega}(\tau_c,\omega_c)\right)^2+\left(2\beta^*\frac{\partial C}{\partial\omega}(\tau_c,\omega_c)\right)^2>0.$$

Consequently, $\mu'(\tau_c) \neq 0$, and the sign of $\mu'(\tau_c)$ is the same as the sign of $-(\partial C/\partial \omega)(\tau_c, \omega_c)$.

In summary, by using (19) and (20), we have obtained that, for $\tau = \tau_c$, Eq. (5) has a pair of simple purely imaginary roots $\pm i\omega_c$ satisfying (18). This completes the proof.

Remark 4.4. If there exists a $k \in \mathbb{N}^*$ such that $\omega_c \tau_c = x_k$, then either $\pm i\omega_c$ are not simple roots of (5) or

$$\left. \frac{\mathrm{d}\,\mathrm{Re}(\lambda)}{\mathrm{d}\tau} \right|_{\tau=\tau_c} = 0.$$

Using a similar argument as in the proof of Proposition 4.2, we obtain that

$$\frac{\partial C}{\partial \omega}(\tau_c,\,\omega_c)=0.$$

Thus, if $\tau_c = -2/(\delta + \beta^*)$, then $\pm i\omega_c$ are not simple roots of (5) and, if $\tau_c \neq -2/(\delta + \beta^*)$, then

$$\frac{\mathrm{d}\operatorname{Re}(\lambda)}{\mathrm{d}\tau}\bigg|_{\tau=\tau_c}=0.$$

In the next theorem, we show that there exists a Hopf bifurcation at the nontrivial equilibrium $x \equiv x^*$ of Eq. (1).

Theorem 4.5. Assume that

$$h(u_0) \leq n \frac{\beta_0 - \delta}{\beta_0} \quad and \quad \delta + \beta^* \neq 0.$$

Then the equilibrium $x = x^*$ is locally asymptotically stable when $0 \le \tau < \tau_0$, where $\tau_0 := \min\{\tau_c; \ \omega_c \tau_c \neq x_k, \ k \in \mathbb{N}\}$, with (τ_c, ω_c) solutions of (15)–(16), and unstable when $\tau_0 \le \tau \le \tau_l$, where τ_l is the largest value of τ_c such that $\omega_c \tau_c \in (x_{2k}, x_{2k+1}), \ k \in \mathbb{N}$. When $\tau = \tau_0$, a Hopf bifurcation occurs at $x \equiv x^*$.

Proof. We first check that $x \equiv x^*$ is locally asymptotically stable when $\tau \in [0, \tau_0)$. Notice that when $\tau \in [0, \tau_0)$, Eq. (5) does not have purely imaginary roots. Let $\tau^* > 0$ be small enough and fixed. Assume that, for $\tau \in (0, \tau^*)$, Eq. (5) has a characteristic root $\lambda(\tau) = \mu(\tau) + i\omega(\tau)$ with $\mu(\tau) > 0$. Separating the real and imaginary parts, we obtain

$$\mu(\tau) = -(\delta + \beta^*) + \frac{2\beta^*}{\tau} \int_0^\tau e^{-\mu(\tau)r} \cos(\omega(\tau)r) dr$$

and

$$\omega(\tau) = -\frac{2\beta^*}{\tau} \int_0^\tau e^{-\mu(\tau)r} \sin(\omega(\tau)r) \,\mathrm{d}r.$$

We deduce that, for $\tau \in (0, \tau^*)$,

$$|\mu(\tau)| \leq |\delta + \beta^*| - 2\beta^*$$
 and $|\omega(\tau)| \leq -2\beta^*$.

Consequently,

$$\lim_{\tau \to 0} \tau \mu(\tau) = 0 \quad \text{and} \quad \lim_{\tau \to 0} \tau \omega(\tau) = 0.$$

Integrating by parts, we obtain

$$2\mu(\tau) = -(\delta + \beta^*) + 2\beta^* \mathrm{e}^{-\tau\mu(\tau)} K(\tau\omega(\tau)).$$

Since $\mu(\tau) > 0$, we have for $\tau \in (0, \tau^*)$ that

$$-(\delta + \beta^*) + 2\beta^* \mathrm{e}^{-\tau\mu(\tau)} K(\tau\omega(\tau)) > 0.$$

When τ tends to zero, we obtain

$$\beta^* - \delta \ge 0.$$

However, $\beta^* - \delta < 0$. This is a contradiction. Therefore, for $\tau \in (0, \tau^*)$, $\mu(\tau) < 0$. Applying Rouché's Theorem [5, p. 248], we obtain that all characteristic roots of (5) have negative real parts when $\tau \in [0, \tau_0)$. Therefore, $x \equiv x^*$ is locally asymptotically stable.

Using Lemma 4.1, Propositions 4.2 and 4.3, we conclude to the existence of τ_l . This concludes the proof. \Box

We illustrate the results of Theorem 4.5 in the next corollary.

Corollary 4.6. Assume that the parameters δ , β_0 and n are given by (2). Then there exists a unique value $\tau_c > 0$ such that the equilibrium $x = x^*$ is locally asymptotically stable when $\tau < \tau_c$ and becomes unstable when $\tau \ge \tau_c$. Moreover, when $\tau = \tau_c$, a Hopf bifurcation occurs at $x \equiv x^*$ and Eq. (1) has a periodic solution with a period close to 46 days (see Fig. 2). The value of τ_c is approximately given by

$$\tau_c \simeq 18$$
 days.

Proof. With the values given by (2), we obtain

$$n \frac{\beta_0 - \delta}{\beta_0} \simeq 2.9153 > h(v_1) \simeq 2.3455.$$

Hence, Proposition 4.2 implies that system (15)–(16) has a unique solution (τ_c, ω_c) with $\tau_c > 0$ and $\omega_c \tau_c \in (0, \pi)$. From Theorem 4.5, we know that the nontrivial equilibrium is locally asymptotically stable when $\tau < \tau_c$, becomes unstable when $\tau \ge \tau_c$, and a Hopf bifurcation occurs at $x \equiv x^*$ for $\tau = \tau_c$. Consequently, for $\tau = \tau_c$, Eq. (1) has a periodic solution with a period close to $2\pi/\omega_c$. One can check that

 $\tau_c \simeq 18$ days and $\omega_c \simeq 0.138$.

Computer simulations confirm our analysis (see Fig. 2). \Box



Fig. 2. With the values given by (2) and $\theta = 1.62 \times 10^8$ cells/kg, Eq. (1) has a periodic solution for $\tau = 18.2$ days. This solution has a period about 50 days (see (a)). The plot of x(t) versus $x(t - \tau)$ (see (b)).

As mentioned earlier, the results in Theorem 4.5 still hold when $\tau_{min} > 0$. However, in this case, the computations in the proof of Theorem 4.5 are much more complicated.

5. Discussion

Hematological diseases have attracted a significant amount of modeling attention because a number of them are periodic in nature [8]. Some of these diseases involve only one blood cell type and are due to the destabilization of peripheral control mechanisms, e.g., periodic auto-immune hemolytic anemia [2,13]. Such periodic hematological diseases involve periods between two and four times the bone marrow production/maturation delay. Other periodic hematological diseases, such as cyclical neutropenia [8], involve oscillations in all of the blood cells and very long period dynamics on the order of weeks to months [6,14] and are thought to be due to a destabilization of the pluripotent stem cell compartment from which all types of mature blood cells are derived.

We have studied a scalar delay model that describes the dynamics of a pluripotent stem cell population involved in the blood production process in the bone marrow. The distributed delay describes the cell cycle duration. We established stability conditions for the model independent of the delay. We have also observed oscillations in the pluripotent stem cell population through Hopf bifurcations. With parameter values given by Mackey [9,10], our calculations indicate that the oscillatory pluripotent stem cell population involves a period of 46 days.

It will be very interesting to study the dynamics of the two-dimensional systems [9,10,12,14] modeling the proliferating phase cells and resting phase cells with distributed delays. We leave this for future consideration.

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